

## AN ANALYSIS OF MICROWAVE DE-EMBEDDING ERRORS

by Lance A. Glasser

Department of Electrical Engineering and Computer Science

and Research Laboratory of Electronics

Massachusetts Institute of Technology  
Cambridge, Massachusetts 02139

**ABSTRACT:** Bounds on de-embedding errors are derived. Both measurement errors and "known" load errors are considered. Of particular interest is the determination of error bounds for unknown loads situated in regions of the Smith chart that are remote from the location of the known de-embedding loads. An example is presented. This analysis is of particular interest for de-embedding microwave diodes.

## INTRODUCTION

In this paper some bounds on de-embedding errors are derived which are especially applicable when measuring microwave diodes. De-embedding is a two-step process. The first step is to unterminate the embedding network  $N$  (see Fig. 1) by the use of known or standard loads<sup>1</sup>. The second step is to use the model that we obtain for the embedding network from the first step, together with measurements of the unknown load taken through the network  $N$ , to estimate the parameters of the unknown load. Frequently we are constrained to use standard loads that are located in only one region of the Smith chart. For instance, varactors, which may be used for de-embedding<sup>2</sup> are limited to the capacitive region of the  $\Gamma_L$ -plane. We are particularly interested here in deriving the error for unknowns that are located in regions of the Smith chart which are remote from the location of the standard loads. For example, given standard loads located at  $\Gamma_L = -1, -j$ , and  $1$ , what is the de-embedding error of an unknown load located at  $\Gamma_L = 0$  or  $\Gamma_L = j$ ? We shall return to this example.

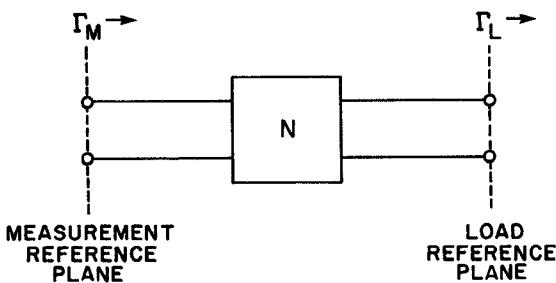


Figure 1 The unknown embedding network  $N$  terminated with a known load  $\Gamma_L$ .

We consider two sources of error.  $\delta_{1M}$ : error in the measurement of the reflection coefficient at the input of network  $N$  and  $\delta_{2L}$ : error in our knowledge of the standard loads.  $\delta_{2L}$  is the radius of a small circle on the  $\Gamma_L$ -plane within which we expect the reflection coefficient of a standard load,  $\Gamma_L$ , to lie.  $\delta_{2L}$  is a function of  $\Gamma_L$ .  $\delta_{1M}$  defines the radius of an uncertainty circle around the measured reflection coefficient  $\Gamma_M$  at the input of network  $N$ .  $\delta_{1M}$  increases with measurement uncertainty.  $\delta_{1M}$  is a function of the input reflection coefficient,  $\Gamma_M$ ; but, since  $\Gamma_M$  is related to  $\Gamma_L$  by a bilinear transformation,  $\delta_{1M}$  is also a function of  $\Gamma_L$ . In this analysis we assume that  $\delta_{1M}$  and  $\delta_{2L}$  have been determined previously either

This work was supported in part by the Joint Services Electronics Program (Contract DAAB07-75-C-1346).

Manuscript received July 21, 1977; revised March 2, 1978.

theoretically or experimentally. For instance, in our work on cryosars<sup>3</sup>,  $\delta_{2L}$  was determined by the measurement scatter among the diode packages that hold MOS capacitor standards. In other words, by using highly redundant data, we could obtain a handle on the size of  $\delta_{2L}$  as a function of  $\Gamma_L$ .  $\delta_{1M}$  may be caused by noise, nonlinearities, or calibration errors in the reflectometer.

## ERROR ANALYSIS

We determine the scattering parameters of network  $N$ , correct to zero order, by the use of a standard unterminating procedure such as that of Bauer and Penfield<sup>1</sup>. From their unterminating calculations we obtain estimates of  $S_{11}$ ,  $S_{22}$ , and  $\Delta$  where  $\Delta \equiv S_{11}S_{22} - S_{12}S_{21}$ . We have

$$\Gamma_M \approx \frac{S_{11} - \Gamma_L \Delta}{1 - \Gamma_L S_{22}}. \quad (1)$$

It is desirable to relate the radius of measurement error,  $\delta_{1M}$ , at the input reference plane to an equivalent radius of error,  $\delta_{2M}$ , at the  $\Gamma_L$  reference plane. We know that such a relation exists because bilinear transformations map circles into circles. We may find this relation by taking the variation of Eq. (1).

$$\frac{d\Gamma_M}{d\Gamma_L} = \frac{S_{11}S_{22} - \Delta}{(1 - \Gamma_L S_{22})^2} \quad (2)$$

or to zero order

$$\delta_{2M} = \left| \frac{(1 - \Gamma_L S_{22})^2}{S_{11}S_{22} - \Delta} \right| \delta_{1M}. \quad (3)$$

We may define a total equivalent error at the output reference plane as

$$\delta'_{2L}(\Gamma_L) \equiv \delta_{2L}(\Gamma_L) + \delta_{2M}(\Gamma_L). \quad (4)$$

Since the data used in the unterminating analysis contain errors, the scattering parameters ( $S_{11}$ ,  $S_{22}$ ,  $\Delta$ ) computed for the network  $N$ , in reality, do not specify  $N$  but rather some network model,  $M$ . An unknown load  $\Gamma_U$  and its corresponding measured input reflection coefficient  $\Gamma_M$  (which we assume for the moment contains no measurement error) are related by the scattering parameters of  $N$ . By invoking the network model  $M$ ,  $\Gamma_U$  is estimated as  $\Gamma'_U$ . Since both  $N$  and  $M$  are perfectly well defined, there exists a bilinear transformation between  $\Gamma'_U$  and  $\Gamma_U$ . We may write this transformation as

$$\Gamma'_U = \frac{A + \Gamma_U}{1 - B - C\Gamma_U}. \quad (5)$$

This is shown schematically in Fig. 2(a).  $M^{-1}$  is the network which undoes the bilinear transformation of network  $M$ . (In MARTHA<sup>4</sup> parlance  $M^{-1} \leftarrow 1$  ZSCALE WN M.) Note that by our construction of the problem the magnitude of the error in Fig. 2(b) is less than  $\delta'_{2L}(\Gamma_L)$ . Because of the assumption that  $S_{11}$ ,  $S_{22}$ , and  $\delta_{2L}(\Gamma_L)$ .

$\Delta$  are correct to zeroth order,  $\Gamma_U' \approx \Gamma_U$  and therefore the magnitudes of A, B, and C are much smaller than one. This may be checked by comparing the measured values of  $\Gamma_M$  with those that can be computed by using (1). They should differ in magnitude by no more than  $\delta_{1M} + \delta_{1L}$  where  $\delta_{1L}$  is defined in a manner analogous to  $\delta_{2M}$  in (3),

$$\delta_{1L} \equiv \left| \frac{s_{11} s_{22} - \Delta}{(1 - \Gamma_L s_{22})^2} \right| \delta_{2L}. \quad (6)$$

The complex de-embedding error  $\Delta_{2U}(\Gamma_U)$ , which includes only errors from the unterminating procedure, is defined as the difference between the estimated and true values of the load reflection coefficient. Expanding (5) to first order, we obtain

$$\Delta_{2U}(\Gamma_U) \equiv \Gamma_U' - \Gamma_U \approx (1, \Gamma_U, \Gamma_U^2) (A, B, C,)^T. \quad (7)$$

We would like to relate the magnitude of  $\Delta_{2U}$  to the equivalent radii of error  $\delta'_{2L}(\Gamma_L)$ . This is done by examining  $\Delta_{2U}(\Gamma_U)$  at three known loads, that is,  $\Delta_{2U}(\Gamma_L)$ . Comparing Figs. 2(a) and 2(b) we observe that for the special case of  $\Gamma_U = \Gamma_L$  (and therefore  $\Gamma_N = \Gamma_M$ ) we may identify  $\Delta_{2U}(\Gamma_L)$  of Fig. 2(a) with the error of Fig. 2(b). We may state

$$|\Delta_{2U}(\Gamma_L)| \leq \delta'_{2L}(\Gamma_L). \quad (8)$$

To obtain the desired bound on  $|\Delta_{2U}(\Gamma_U)|$  we first evaluate (7) at  $\Gamma_U = \Gamma_L$ . A relation between the parameters A, B, and C, and  $\Delta_{2U}(\Gamma_L)$  is found for the three loads under consideration. A bound on  $\Delta_{2U}(\Gamma_U)$  may then be found by invoking the triangle inequality.

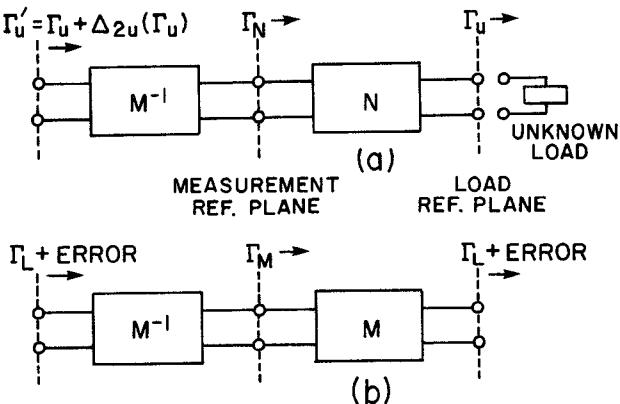


Figure 2 (a) Schematic of how the network model M is used to de-embed an unknown load  $\Gamma_U$ . The resulting error is  $\Delta_{2U}(\Gamma_U)$ . (b) A useful network identity from the definition of M. The magnitude of the error is less than  $\delta'_{2L}$  by definition.

#### AN EXAMPLE

The procedure is best illustrated by an example. Consider, for instance, the three loads  $\Gamma_L = 1, -1, -j$ . Evaluating (7) for these loads and letting  $\alpha_1 \equiv \Delta_{2U}(\Gamma_L = 1)$ ,  $\alpha_2 \equiv \Delta_{2U}(\Gamma_L = -1)$ , and  $\alpha_3 \equiv \Delta_{2U}(\Gamma_L = -j)$ , we obtain

$$(\alpha_1, \alpha_2, \alpha_3) = (A, B, C) H^T$$

where

$$H \equiv \begin{pmatrix} 1 & 1 & 1 \\ 1 & -1 & 1 \\ 1 & -j & -1 \end{pmatrix}. \quad (9)$$

$H$  is a function of the three standard loads. Combining (7) and (8), we get

$$\Delta_{2U} = (1, \Gamma_U, \Gamma_U^2) H^{-1} (\alpha_1, \alpha_2, \alpha_3)^T \quad (10)$$

where  $H^{-1}$  may be evaluated as

$$H^{-1} = \frac{1}{4} \begin{pmatrix} 1+j & 1-j & 2 \\ 2 & -2 & 0 \\ 1-j & 1+j & -2 \end{pmatrix}. \quad (11)$$

$\alpha_1$ ,  $\alpha_2$ , and  $\alpha_3$ , however, represent the equivalent error of the standard loads that are known in terms of  $\delta'_{2L}$  from the first part of this article. Therefore we write

$$|\alpha_1| \leq \delta'_{2L}(\Gamma_L = 1); |\alpha_2| \leq \delta'_{2L}(\Gamma_L = -1);$$

$$|\alpha_3| \leq \delta'_{2L}(\Gamma_L = -j). \quad (12)$$

By using the triangle inequality we can rewrite (10)

$$\begin{aligned} |\Delta_{2U}(\Gamma_U)| &\leq \frac{1}{4} |1 + j + 2\Gamma_U + \Gamma_U^2 - j\Gamma_U^2| \delta'_{2L}(\Gamma_L = 1) \\ &\quad + \frac{1}{4} |1 - j - 2\Gamma_U + \Gamma_U^2 + j\Gamma_U^2| \delta'_{2L}(\Gamma_L = -1) \\ &\quad + \frac{1}{4} |2 - 2\Gamma_U^2| \delta'_{2L}(\Gamma_L = -j). \end{aligned} \quad (13)$$

It is interesting to evaluate (13) for the two unknown loads mentioned earlier. We also assume  $\delta'_{2L}(\Gamma_L = 1) = \delta'_{2L}(\Gamma_L = -j) = \delta'_{2L}(\Gamma_L = -1) \equiv \varepsilon$ . Under these conditions,

$$\begin{aligned} |\Delta_{2U}(\Gamma_U = 0)| &\leq 1.2\varepsilon \\ |\Delta_{2U}(\Gamma_U = j)| &\leq 3\varepsilon. \end{aligned} \quad (14)$$

This illustrates clearly how the error grows in regions of the Smith chart that are remote from the location of the standard loads. The total radius of error for  $\Gamma_U$  at the load reference plane is

$$\delta_{\text{TOTAL}}(\Gamma_U) = |\Delta_{2U}(\Gamma_U)| + \delta_{2M}. \quad (15)$$

Although it is demonstrated by a specific example, the error analysis outlined here is very general. It can also be programmed easily for computer or calculator use.

#### ACKNOWLEDGEMENT

The author wishes to thank R. L. Kyhl for his guidance and suggestions, and P. Penfield, Jr., M. S. Gupta, R. Bauer, and J. F. Cooper for their helpful discussions.

#### REFERENCES

1. R. F. Bauer and P. Penfield, Jr., "De-Embedding and Unterminating," *IEEE Trans. Microwave Theory and Techniques*, vol. MTT-22, pp. 282-288, March, 1974.
2. C. E. Simmons, "A Cooled K-Band Mixer Using Schottky Barrier Diodes." S. M. Thesis, Department of Electrical Engineering and Computer Science, MIT, Cambridge, MA., December, 1974.
3. L. A. Glasser and R. L. Kyhl, "The Silicon Cryosar at Microwave Frequencies", unpublished.
4. P. Penfield, Jr., *MARTHA USERS MANUAL*, Cambridge, MA.: MIT Press, 1971.